

Carbon nanotubes as ultra-high quality factor mechanical resonators — and much more!

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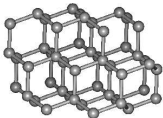
Institute for Experimental and Applied Physics,
Universität Regensburg, Germany

Condensed Matter and Materials Physics (CMMP10)
University of Warwick, Coventry, UK

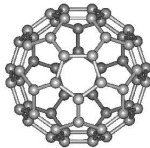
14 December 2010

Carbon nanotubes: a more exciting

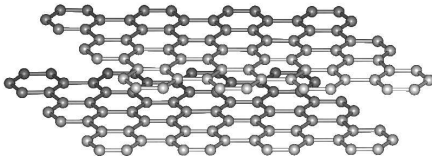
(and not so flat) form of carbon



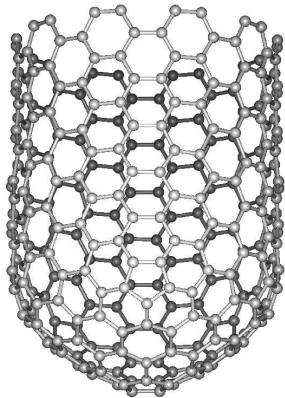
diamond



fullerene (C₆₀)



graphite / graphene



nanotube

Mechanical properties of carbon nanotubes

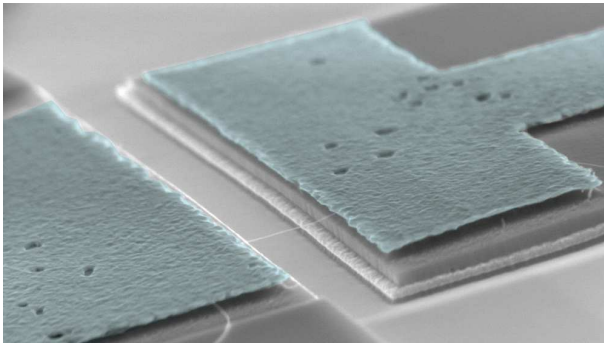
- stiffer than steel
- resistant to damage from physical forces
- very light

- Young's modulus $E = \frac{F/A}{\Delta L/L}$:
 $E_{\text{CNT}} \simeq 1.2 \text{TPa}$, $E_{\text{steel}} \simeq 0.2 \text{TPa}$
- Density:
 $\rho_{\text{CNT}} \simeq 1.3 \frac{\text{g}}{\text{cm}^3}$, $\rho_{\text{Al}} \simeq 2.7 \frac{\text{g}}{\text{cm}^3}$

- (still) “material of dreams”

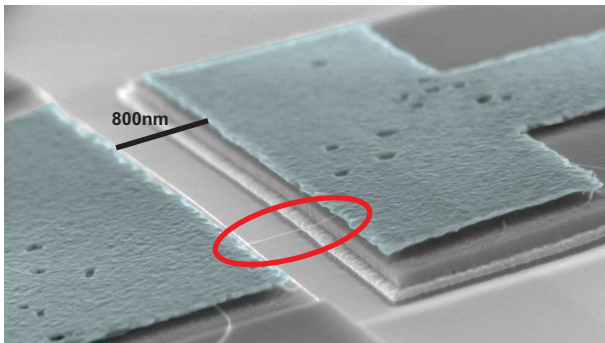


Doubly clamped nanotube resonators



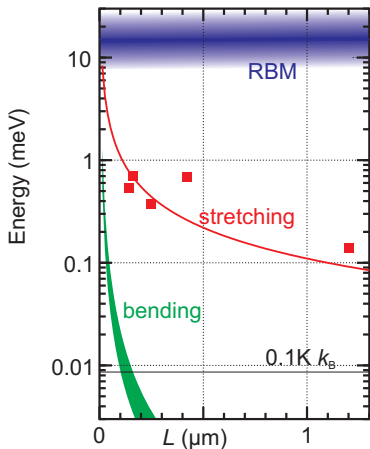
nanotube is suspended like a guitar or violin string
low mass, high stiffness \rightarrow high resonance frequency, large quantum effects
single clean macromolecule \rightarrow low dissipation???

Doubly clamped nanotube resonators



nanotube is suspended like a guitar or violin string
low mass, high stiffness \rightarrow high resonance frequency, large quantum effects
single clean macromolecule \rightarrow low dissipation???

Vibration modes of carbon nanotubes



- **stretching** (longitudinal) mode:

$$h\nu \propto L^{-1}$$

$$h\nu = 1100 \dots 110 \mu\text{eV},$$

$$\nu = 270 \dots 27 \text{GHz}$$

(for 100nm...1 μm)

- **bending** (transversal) mode:

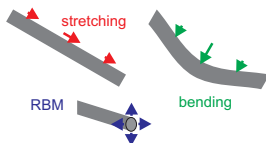
$$h\nu \propto L^{-2}$$

$$h\nu = 10 \dots 0.1 \mu\text{eV},$$

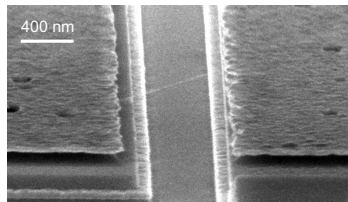
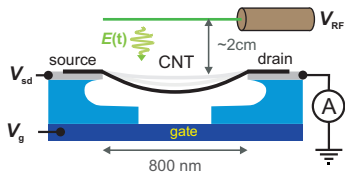
$$\nu = 2.4 \text{GHz} \dots 24 \text{MHz}$$

(for 100nm...1 μm)

$h\nu \propto d$, also tension-dependent



Chip fabrication and measurement setup

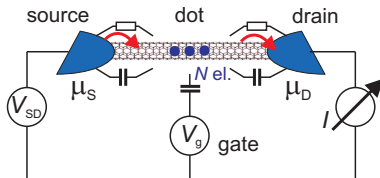


- First make chip (Pt electrodes, trench)
- Then CVD-grow nanotubes across electrodes
- Back gate connected to a gate voltage source V_g
- RF antenna suspended ~ 2 cm above chip
- Dilution refrigerator ($T \simeq 20$ mK)
- **Only dc measurement**

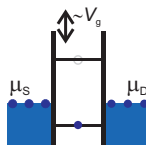
Low-temperature transport: Coulomb blockade

dilution refrigerator $T \lesssim 20\text{ mK}$

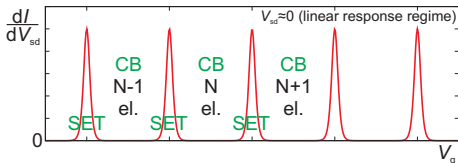
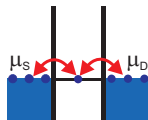
- Tunnel barriers between leads and nanotube
- Low temperature $k_B T \ll e^2/C$: formation of a **quantum dot**



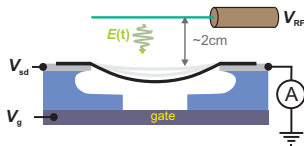
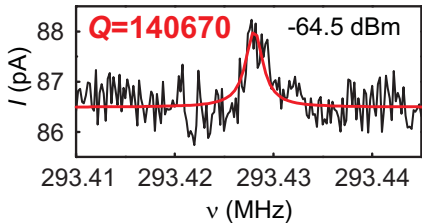
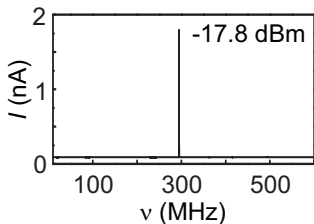
Coulomb blockade



Single electron tunneling

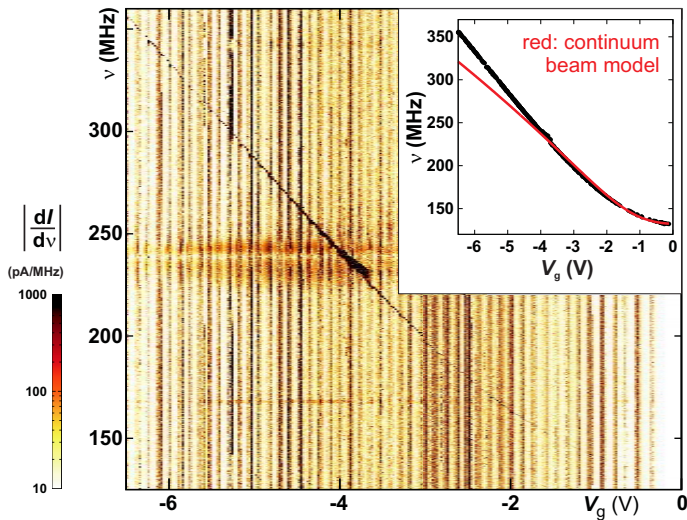


Fixed V_g and V_{SD} , sweep of RF signal frequency



- Sharp resonant structure in $I_{dc}(\nu)$
- Very low driving power required
- High $Q = \nu/\Delta\nu$ ($\Delta\nu = \text{FWHM}$)

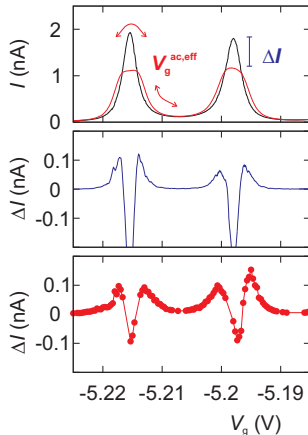
V_g dependence — this is really a mechanical resonance!



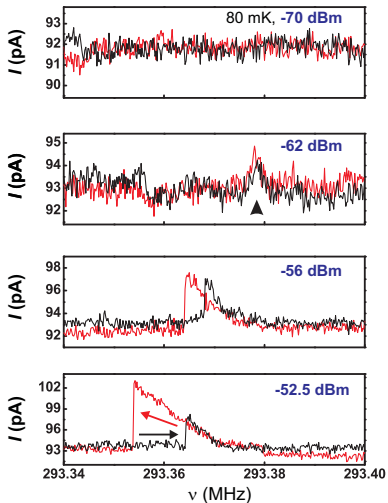
larger $|V_g| \rightarrow$ increased tension \rightarrow higher frequency ν

Detection mechanism — mechanically induced averaging

- at resonant driving the nanotube position oscillates
- oscillating C_g
→ fast averaging over $I(V_g)$
- black line: dc measurement $I(V_g)$
- red line: this numerically averaged
- blue: difference, effect of averaging
- red points: measured peak amplitude in $I(v)$, for different values of V_g

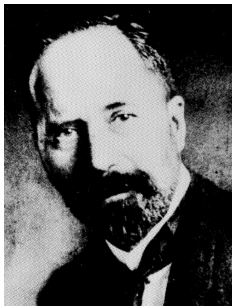


Driving into nonlinear response...



- same temperature
- same working point V_g, V_{SD}
- low driving power:
symmetric, “linear” response
- high driving power:
asymmetric response, hysteresis
Duffing-like oscillator

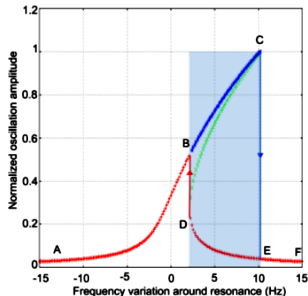
Georg Duffing (1861 – 1944) and his oscillator



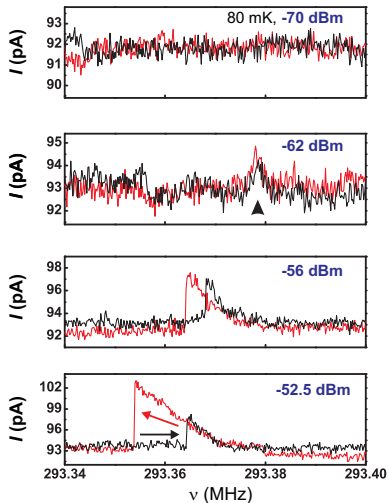
Duffing differential equation:

$$m\ddot{x} + cx + bx^3 = F \sin \omega t$$

- Driven mechanical oscillator with non-linear response
- Response becomes bistable
→ large or small amplitude
- Switching between branches

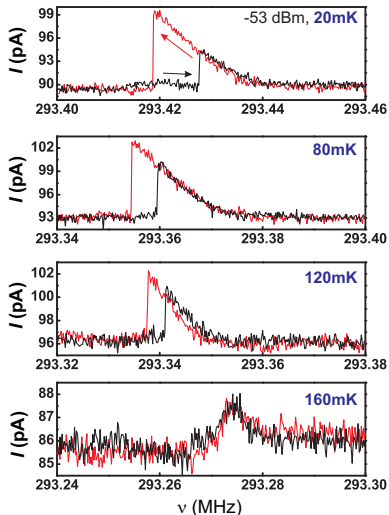


Driving into nonlinear response...



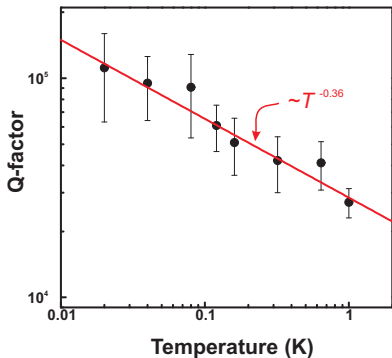
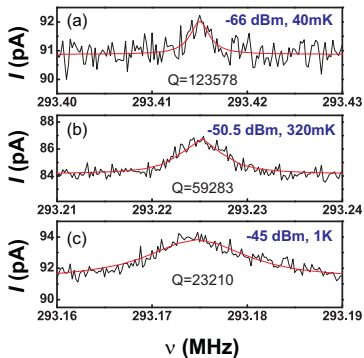
- same temperature
- same working point V_g, V_{SD}
- low driving power:
symmetric, “linear” response
- high driving power:
asymmetric response, hysteresis
Duffing-like oscillator

... and then increasing the temperature



- same driving power
- same working point V_G, V_{SD}
- low temperature:
asymmetric response, hysteresis
Duffing-like oscillator
- high temperature:
symmetric, "linear" response
peak broadening

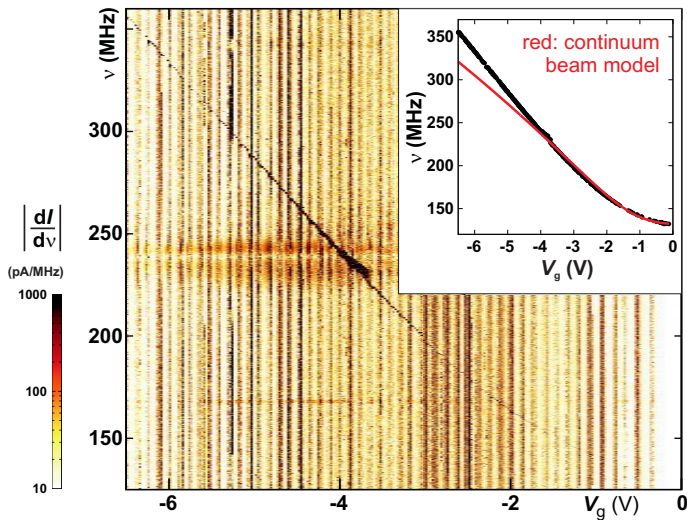
Temperature dependence of Q



$Q(T)$ fits power law prediction for intrinsic dissipation in nanotube

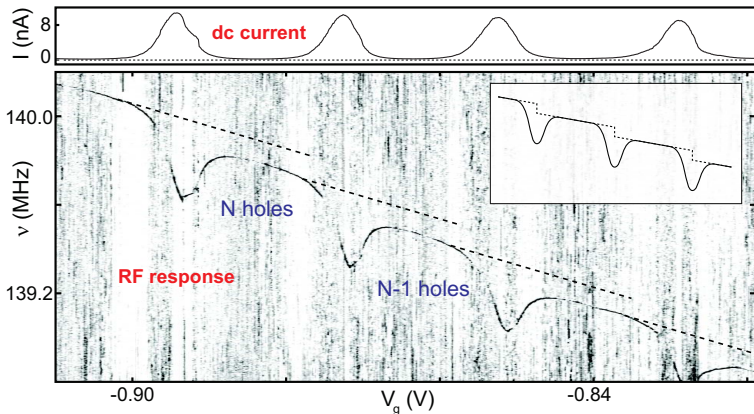
→ H. Jiang *et al.*, Phys. Rev. Lett. **93**, 185501 (2004)

V_g dependence — this is really a mechanical resonance!



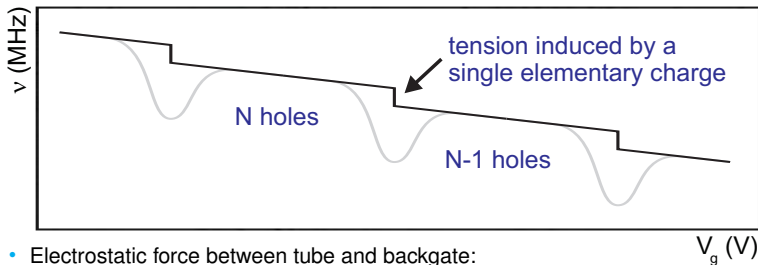
larger $|V_g| \rightarrow$ increased tension \rightarrow higher frequency ν

Detailed $\nu(V_g)$: with current, frequency decreases



“Coulomb blockade oscillations of **mechanical resonance frequency**”
electrostatic contribution to spring constant

Model for $\nu(V_g)$ – part I: “slope and steps”



- Electrostatic force between tube and backgate:

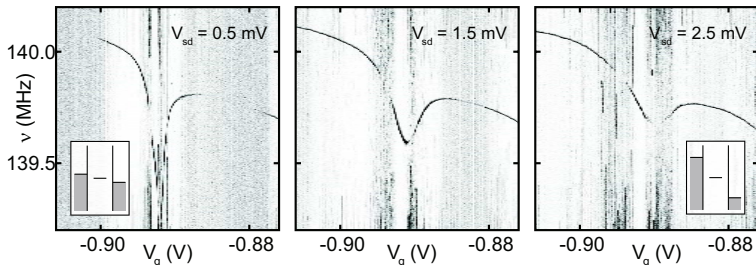
$$F_{\text{dot}} = \frac{1}{2} \frac{dC_g}{dz} (V_g - V_{\text{dot}})^2$$

- Quantum dot voltage:

$$V_{\text{dot}} = \frac{C_g V_g + q_{\text{dot}}}{C_{\text{dot}}}, \quad q_{\text{dot}}(q_c) = -|e| \langle N \rangle (q_c), \quad q_c = C_g V_g$$

- Overall slope: continuous increase of voltage V_g on gate
- Steps: discrete change of V_{dot} (single elementary charges!)

Model for $\nu(V_g)$ – part II: “steps become dips”

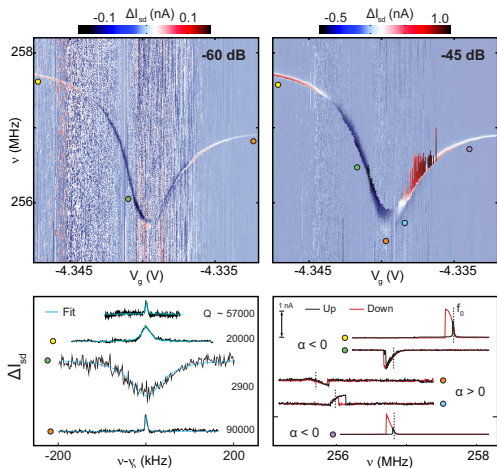


- $q_c = C_g(z)V_g$ is function of z
- Electrostatic contribution to spring constant:

$$k_{\text{dot}} = -\frac{dF_{\text{dot}}}{dz} = \frac{V_g(V_g - V_{\text{dot}})}{C_{\text{dot}}} \left(\frac{dC_g}{dz} \right)^2 \left(1 - |e| \frac{d\langle N \rangle}{dq_c} \right)$$

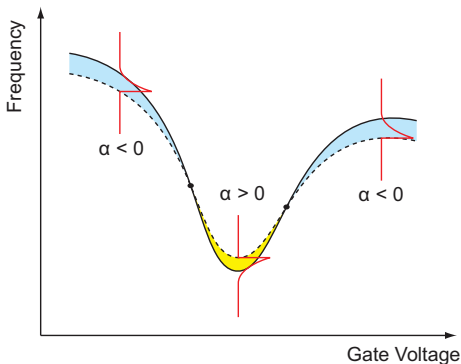
- Always negative, always decreasing frequency

Also mechanical Q and nonlinearity dominated by current



- Dissipation whenever charge can fluctuate
- Q decreases on SET peaks
- Nonlinearity dominated by tunneling
- Switches between weakening and softening spring

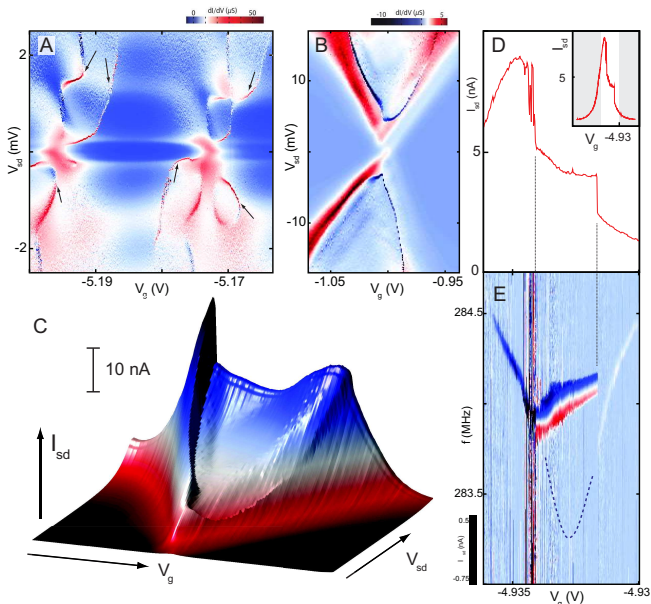
Interaction-induced nonlinearity $\alpha(V_g)$



$$\alpha_{\text{dot}} = -\frac{d^3 F}{dz^3} = \frac{d^2}{dz^2} k_{\text{dot}}(q_c) = V_g^2 \left(\frac{dC_g}{dz} \right)^2 \frac{d^2 k_{\text{dot}}}{dq_c^2}$$

The sign of α_{dot} follows the sign of the curvature of k_{dot} .

Self-excitation of the resonator

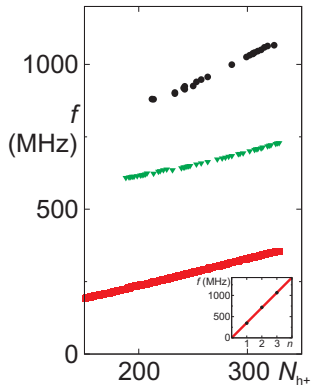


What do we have so far?

- Mechanical resonator, $120 \text{ MHz} \lesssim \nu \lesssim 360 \text{ MHz}$, $Q \lesssim 150000$
- Easy driving into nonlinear oscillator regime
- Single-electron steps of the resonance frequency
- Backaction of single electron tunneling on ν , Q , nonlinearity

- Estimated motion amplitude at resonant driving $\sim 250 \text{ pm}$
compare thermal motion 6.5 pm , zero-point motion 1.9 pm
- Application as mass sensor: **sensitivity** $4.2 \frac{u}{\sqrt{\text{Hz}}}$
- Without driving: **mechanical thermal occupation** $n \simeq 1.2$

Higher frequency (I): higher vibration modes

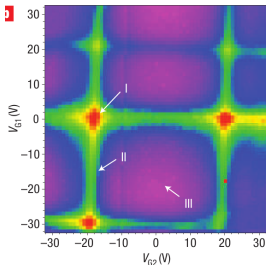
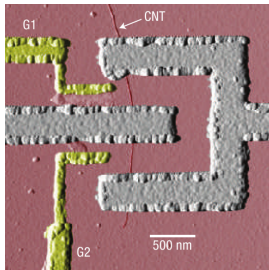


- higher harmonics visible too
- dc current signal is smaller (node(s) in nanotube motion, smaller change in total capacitance)
- at high tension, integer frequency multiples (expected for a string resonator)

Higher frequency (II): just make it shorter!

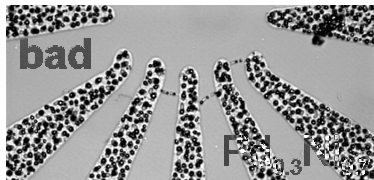
:) ongoing work in Delft and Regensburg :)

Going super



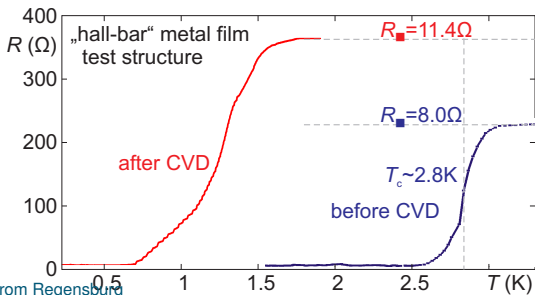
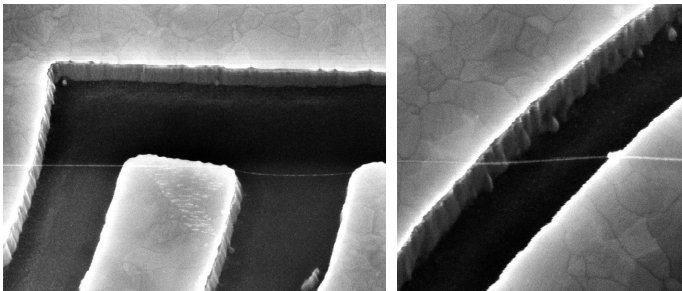
- Nanotubes can carry supercurrents via proximity effect
 - Use superconducting electrodes
 - Cooper pair tunneling
 - Nanotube SQUIDs, ac Josephson effect, intrinsic cooling of the vibration, ...
-
- image: example for beautiful (non-suspended) hybrid device
 - Superconducting support and control electronics!

Pitfalls and problems for *ultra-clean samples*



- need to first prepare on-chip infrastructure: contacts, gates, trenches, ...
- then grow nanotubes across the chip with **CVD as last step**
- 10min, 900°C, CH₄ and H₂: for a metal thin film “as bad as it gets”
- melting, recrystallization
 - deformation, loss of conductivity
- hydrogen / carbon storage in metal
 - lowering of superconductor T_c
- influence of metal on nanotube growth?
- properties of nanotube–metal contact?

but... it seems to be working

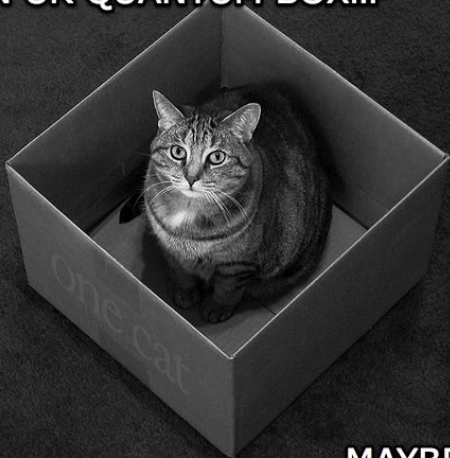


lots of opportunities

- beam resonator in quantum mechanical ground state
- transition classical – quantum harmonic oscillator
- quantum nonlinear resonator properties (many theory predictions!)
- ...
- ...
- ...

Go quantum limit!

IN UR QUANTUM BOX...



...MAYBE.

The old team at TU Delft

Thanks!



Gary Steele



Benoit Witkamp



Menno Poot



Leo Kouwenhoven



Harold Meerwaldt



Herre van der Zant

My **new** team at Uni Regensburg

Thanks!



Daniel Schmid



Dominik Preusche



Peter Stiller



you?



Christoph Strunk

and everyone else!