Co-sputtered MoRe thin films for carbon nanotube growth-compatible superconducting coplanar resonators

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Molybdenum rhenium alloy thin films can exhibit superconductivity up to critical temperatures of $T_c = 15$ K. At the same time, the films are highly stable in the high-temperature methane / hydrogen atmosphere typically required to grow single wall carbon nanotubes. We characterize molybdenum rhenium alloy films deposited via simultaneous sputtering from two sources, with respect to their composition as function of sputter parameters and their electronic dc as well as GHz properties at low temperature. Specific emphasis is placed on the effect of the carbon nanotube growth conditions on the film. Superconducting coplanar waveguide resonators are defined lithographically; we demonstrate that the resonators remain functional when undergoing nanotube growth conditions, and characterize their properties as function of temperature. This paves the way for ultra-clean nanotube devices grown in situ onto superconducting coplanar waveguide circuit elements.

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I. INTRODUCTION

The discovery of single-walled carbon nanotubes [1, 2] opened the door for a wide range of both fundamental research and technical applications in such electronic and nano-electromechanical systems. These macromolecules exhibit particular properties as, e.g., high mechanical, chemical and thermal stability as well as excellent electronic and thermal conductance. They can carry high currents, display outstanding room temperature transistor characteristics, and may compete in the future with or even replace conventional silicon-based logic devices [3–5]. From a fundamental point of view carbon nanotubes provide a highly promising material for quantum information and computation. The macromolecules act as quasi one-dimensional conductors because of their large aspect ratio; the electrostatic definition of localized potential wells can trap electrons in a nuclear spin free environment with well-defined quantum levels [6–10].

One way to ensure that the carbon nanotubes remain clean and defect-free is to grow them directly on pre-fabricated contact electrodes, such that subsequently no further lithography and / or wet chemistry takes place. This technique has led to a multitude of striking measurement results, from single quantum dot spectroscopy [9, 11] to the characterization of high-$Q$ vibrational modes [12, 13]. Performing high-frequency measurements on single wall carbon nanotube structures is a logical extension. Many techniques applied for such measurements on mesoscopic structures require the use of on-chip radiofrequency electronics, in particular superconducting coplanar waveguides and microwave resonators formed by these [15, 16]. An example for such a possible combined device structure is sketched in Fig. 1(a) – a carbon nanotube is capacitively coupled to the coplanar $\lambda/2$ resonator. In combination with the overgrowth process, this poses extensive difficulties, since the carbon nanotube growth process — typically 10-15 min in a methane–hydrogen atmosphere [17] — is highly detrimental to superconducting thin films. The molybdenum target is driven by a radiofrequency source with impedance matching, the rhenium target using a dc source. A mass flow controller sets the argon flow into chamber.

![FIG. 1. (Color online) (a) Schematic sketch of a carbon nanotube quantum dot device capacitively coupled to a coplanar $\lambda/2$ resonator. For similar devices using different material systems, see Refs. 14 and 15. (b) Sketch of the Orion dual-source sputtering setup used during fabrication of our metal thin films. The molybdenum target is driven by a radiofrequency source with impedance matching, the rhenium target using a dc source. A mass flow controller sets the argon flow into chamber.](image)

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CVD growth conditions and subsequently still exhibit superconducting behavior [23]. In addition, after in-situ growth they provide transparent electronic contacts to carbon nanotubes [9, 23, 24].

In the following, we characterize molybdenum rhenium alloy films deposited via simultaneous sputtering from two sources, i.e., co-sputtering, with respect to their composition as function of sputter parameters and their dc as well as GHz properties at low temperature. Specific emphasis is placed on the effect of the carbon nanotube growth conditions on the film. Coplanar waveguide λ/4 resonators are defined lithographically and characterized at dilution refrigerator temperatures; even after undergoing the growth conditions internal quality factors of up to \( Q_i \approx 5000 \) can be found. The temperature dependence of the resonance frequency and the internal quality factor is evaluated and found consistent with theoretical models.

II. THIN FILM DEPOSITION

Figure 1(b) sketches the setup used for the film fabrication. The UHV-chamber with typical pressures in the range of \( 10^{-8} \) to \( 10^{-7} \) mbar contains two sputter targets. The sample holder is placed approximately 13.5 cm above and in the middle of both targets which are mounted at a distance of circa 15 cm to each other. Argon gas is injected via a mass flow controller close to the molybdenum target, and a plasma is ignited using a radiofrequency drive at a chamber pressure of \( 10^{-1} \) mbar.

Subsequently the chamber pressure is reduced to \( 7 \cdot 10^{-3} \) mbar and the argon plasma at the rhenium electrode is ignited using a dc power supply; at the same time the rf-output is adjusted such that the plasma close to the molybdenum electrode and target remains stable. As soon as both targets are sputtered, the shutters are opened and the deposition of the MoRe alloy starts. By keeping the chamber pressure constant over the whole process time, the deposition rate of both materials is kept approximately constant. In the following the alloy composition of the films is varied only by tuning one parameter, namely the output power \( P_{\text{Mo}} \) of the rf-generator at the molybdenum target.

III. RESULTING ALLOY

In order to obtain the alloy composition, X-ray photoelectron spectroscopy (XPS) is performed on the cosputtered films [25, 26]. Within a predefined area of four to ten square millimeters the samples are irradiated by a monochromatic X-ray source, and the resulting emitted photoelectrons are collected and spectroscopically analyzed with respect to energy and intensity. The chemical sensitivity is given by the element-specific distribution of binding energies \( E_B \) of the electronic core levels.

Characteristic XPS data of a cosputtered film are plotted in Fig. 2(a). The signal peaks corresponding to atomic and molecular core levels have been identified following Ref. 27. The black curve in Fig. 2(a) reveals the chemical composition of the as-grown sample surface, the red line was recorded after removing approximately 5 nm of the film by in-situ argon-ion sputtering beam sputtering within the XPS chamber. Normalized detail spectra close to the Mo 3d-peaks (b) and the Re 4f-peaks (c) (see text). The energetic step size is set to 0.2 eV. (d) Depth profile of the atomic composition, calculated from XPS spectra as in (a), from surface to ~5 nm depth (subsequently named bulk). The resulting MoRe-alloy composition after each step is stated above. (e) Relative molybdenum atomic percentage in the bulk film for different power settings \( P_{\text{Mo}} \) at the molybdenum target (see text).
and carbon peaks are observed. This is likely due to the fact that all samples have been exposed to air during transfer from the sputtering device into the UHV chamber of the XPS setup.

Adsorbates and the film itself can be etched by in situ Ar-sputtering for obtaining spectra from lower layers. An etching step lasts several minutes at a chamber pressure of less than 3·10^{-8} mbar in which the base pressure is of 5 · 10^{-10} mbar. Etching steps are repeated until the oxygen 1s-peak only negligibly contributes to the whole spectrum (at most about 10%). The red (gray) line in Fig. 2(a) displays the corresponding spectrum. Note that also the carbon peak is now strongly suppressed. Subsequent tests using a profilometer show that approximately 5 nm of the films were removed. All spectra obtained after such a corresponding etching step are henceforth denoted as “bulk” spectra.

Figure 2(b) and 2(c) display details of the spectrum close to the Mo 3d- and Re 4f- peaks, referenced to the C 1s-peak.[26] The additional structure in the Mo 3d surface spectrum originates from MoO_3 and MoO_2 forming at the alloy surface [25, 28]. It decreases subsequently until disappearing in the bulk where only the two Mo 3d_{5/2} and Mo 3d_{3/2} peaks are part of the spectrum. In contrast, at the Re 4f peak neither changes in the line shape nor peak shifts are observed in the surface compared to the bulk spectrum, indicating the absence of rhenium oxides at the surface and in the bulk layers. All films examined in this work exhibit the same behavior for the rhenium peaks [25, 26].

All rhenium peaks of the film are shifted by 1.3 eV to higher \(E_B\)-values, compared to the literature values of pure rhenium. Similar values have been identified in Ref. 26 as chemical shift due to the Mo-Re compound formation.

Using the method of area sensitivity factors and evaluating the Mo 3d-, Re 4f-, C 1s-, and O 1s-peaks, the atomic concentrations of the sample have been estimated [27]. Fig. 2(d) displays the atomic concentrations as a function of depth. While the exposure to air leads to significant carbon and oxygen percentages at the surfaces, these both decrease strongly; for the “bulk” spectrum about 10% of carbon remains.

Subsequently the molybdenum-rhenium alloy ratio is obtained by normalizing to the sum of both sputtered metals. As can be clearly seen in Fig. 2(e), where the resulting alloy ratio in the bulk film is plotted as a function of rf power \(P_{Mo}\) applied to the molybdenum sputter source, the resulting molybdenum contribution in the alloy can be controlled over a wide range. The solid line in Fig. 2(e) is a quadratic fit to the data points.

IV. INFLUENCE OF THE NANOTUBE GROWTH ENVIRONMENT

Simulating the carbon nanotube chemical vapor deposition (CVD) growth process typically used to locally grow few clean single-wall carbon nanotubes [17], the chip including the MoRe thin film is heated up in an argon and hydrogen gas flow, then exposed to a methane-hydrogen atmosphere at about 900 °C for several minutes and subsequently cooled down under argon and hydrogen flow.

The XPS of a molybdenum rhenium alloy sputtered with \(P_{Mo} = 75\) W after exposure to the CVD environment is shown in Fig. 3(a), and a normalized detail plot of the Mo 3d peak in Fig. 3(b). A strong carbon peak is visible even in the bulk. Furthermore in both surface and bulk molybdenum spectra oxide peaks are not observed, as visible in Fig. 2(b).

Having in mind that molybdenum oxides exhibit drastically lower melting points of 795 °C for MoO_3, and 1100 °C for MoO_2 [29] than pure molybdenum with 2610 °C [30], the data indicates that reduction of the oxide to atomic molybdenum takes place during CVD.

An additionally possible process is the formation of molybdenum carbides. The expected XPS peak of Mo_2C \((E_B = 227.75\) eV) is very close to that of Mo 3d_{5/2} \((E_B = 228\) eV), which makes detection challenging with our experimental resolution. Interestingly, molybdenum carbides display superconductivity with critical temperatures \(6\) K ≤ \(T_c\) ≤ 9 K [31]. However, typically the growth of molybdenum carbide out of pure molybdenum, for example in a nitrogen-xylene atmosphere, takes place at temperatures higher than 900 °C, see [31].
by means of optical lithography and SF$_6$/Ar reactive ion etching (see Fig. 4(a)). Afterwards, selected structures are placed into the CVD furnace and exposed to the nanotube growth environment for several minutes. All measurements have been performed on films either examined by XPS (see last two sections) or deposited simultaneously to these in the same deposition step. Devices using two different alloy compositions have been examined, namely Mo$_{20}$Re$_{80}$ and Mo$_{58}$Re$_{42}$ obtained with $P_{H_2}$ = 75 W and $P_{CH_4}$ = 200 W, respectively.

At room temperature, compared to Mo$_{58}$Re$_{42}$ both resistivity and sheet resistance of the Mo$_{20}$Re$_{80}$ samples are higher by a factor 3–5: before CVD we obtain $\rho \approx 3.0 \cdot 10^{-7}$ $\Omega$ m for Mo$_{58}$Re$_{42}$ and $\rho \approx 9.0 \cdot 10^{-7}$ $\Omega$ m for Mo$_{20}$Re$_{80}$. Resistances slightly increase during exposure to the CVD environment, to $\rho \approx 4.0 \cdot 10^{-7}$ $\Omega$ m for Mo$_{58}$Re$_{42}$ and $13 \cdot 10^{-7}$ $\Omega$ m for Mo$_{20}$Re$_{80}$.

Results of low-temperature measurements performed on Mo$_{20}$Re$_{80}$ devices are plotted in Fig. 4(b) and Fig. 4(c). Independent from CVD-exposure, the residual-resistance values $RRR$ for all Mo$_{20}$Re$_{80}$ devices are in the range 0.8 $\leq RRR \leq 1.0$. Fig. 4(b) displays the critical current of a Mo$_{20}$Re$_{80}$ film (Hall bar geometry B, with $T_c = 9.2$ K) after 10 min CVD exposure. It carries a supercurrent up to $I_c \geq 80$ mA, corresponding to a critical current density of $j_c \geq 2.7 \cdot 10^3$ A/mm$^2$ over a wide temperature range up to ca. 5 K before the transition to a normal conductor takes place above 8 K. This very high value for $j_c$ is well in accordance with the results of [32], where after high-temperature annealing critical current densities of up to $1.8 \cdot 10^5$ A/mm$^2$ through suspended Mo$_{58}$Re$_{42}$-nanostructures were reported.

A second Mo$_{20}$Re$_{80}$ Hall bar device (type A, $T_c = 8.3$ K), not exposed to CVD, exhibits a critical current $I_c \approx 114$ mA at $T = 4.2$ K, corresponding to a lower current density $j_c = 2.8 \cdot 10^3$ A/mm$^2$. Also much longer exposure to the CVD environment again lowers the reachable critical current density.

Figure 4(c) displays data on the magnetic field dependence of the critical temperature. Fitting the empirical relation $B_c(T) = B_{c0} \cdot \left(1 - (T/T_{c0})^2 \right)$ [33] results in high characteristic values $B_{c0} = 7.3$ T and $T_{c0} = 8.3$ K expected for molybdenum rhenium alloys, where $T_{c0}$ denotes the zero field critical temperature and $B_{c0}$ the extrapolated critical field at zero temperature. Again, as observed for critical currents, prolonged (30 min) CVD exposure results in a strong decrease of both values to here $B_{c0} = 3.6$ T and $T_{c0} = 3.9$ K [32, 34]. This effect is also visible in Figure 4(d), displaying the critical temperature for two different alloy compositions. Even for only 10 minutes growth time, $T_c$ of the Mo$_{58}$Re$_{42}$ film decreases to the half, well in accordance with the results of Ref. 34. Interestingly, in contrast the Mo$_{20}$Re$_{80}$ alloy keeps its critical temperature range of $8 K \leq T_c \leq 9 K$ for growth times up to 20 minutes, only reaching $< 4 K$ after 30 minutes of exposure to the CH$_4$/H$_2$ flow.

V. DC CHARACTERIZATION

To characterize the electronic properties of the co-sputtered films at room temperature as well as at cryogenic temperatures, the fabricated thin films on top of SiO$_2$ or Al$_2$O$_3$ substrates are patterned into Hall bars

![Hall bar device](image)

FIG. 4. (Color online) (a) Exemplary SEM micrograph of a test Hall bar structure, type A. Dimensions are for type A width $W_A = 28$ $\mu$m and film thickness $d_A \approx 150$ nm, for type B $W_B = 5$ $\mu$m and $d_B \approx 60$ nm. (b) Measured critical current $I_c$ through a Hall bar device (type B), Mo$_{20}$Re$_{80}$, as a function of temperature. (c) Critical temperature $T_c$ of Mo$_{20}$Re$_{80}$ films as a function of applied magnetic field $B$, with and without CVD growth exposure. (d) Measured critical temperature $T_c$ as a function of CH$_4$/H$_2$ flow time during the CVD process. Straight lines are guides to the eye.

The area sensitivity factor analyzing method results in a bulk composition of 26% carbon, 22% molybdenum, and 49% rhenium plus neglectable oxygen residues. Since the etching time for the bulk spectra has been kept constant, it is obvious that the high carbon contribution does not have its origin in atmospheric adsorbates but in diffusion of carbon into the alloy during CVD. The relative atomic ratio of the sputtered metals is Mo$_{31}$Re$_{69}$, indicating structural changes during CVD.

Spectroscopy on a second sample using twice the gas flow of methane results in a composition of 12% molybdenum, 23% rhenium, and even 58% carbon in the bulk. From this we conclude that the penetration of carbon into the alloy also increases. However, the resulting sputtered metal ratio of Mo$_{34}$Re$_{66}$ still remains close to the previous sample.
VI. COPLANAR RESONATOR DEVICES

Coplanar waveguide λ/4 resonators were fabricated to investigate the high frequency behavior of the alloy material and its suitability for cavity quantum electrodynamics and optomechanics experiments. After sputtering of 150 nm MoRe on SiO₂ or Al₂O₃ substrates, the structures were patterned using optical lithography and reactive ion etching with SF₆/Ar. A micrograph of a device coupling three λ/4 resonators to a common feed line is shown in Fig. 5(a). The devices are glued on a printed circuit board, bonded with aluminum bond wires and subsequently characterized in a dilution refrigerator. On the signal input side the experimental wiring of superconducting UTF85 semirigid NbTi cables includes attenuators as thermal anchoring at every temperature stage. The input signal is attenuated by approximately 53 dB, transmitted through the device under test and then amplified by 29 dB by a low noise HEMT amplifier[35] at the 1 K stage.

Near its fundamental resonance frequency, given by the length of the resonator l and an effective permittivity εeff,

$$f_r = \frac{1}{\sqrt{\varepsilon_{eff}}} \frac{c}{4l},$$

each of the three resonators behaves like a parallel lumped-element RLC circuit, coupling energy out of the feed line and leading to a distinct resonant drop in feedline transmission S₂₁. In the vicinity of the resonance, the transmitted signal can be expressed by [36]

$$S_{21} = 1 - \frac{Q_i}{Q_i} e^{\frac{\pi}{f_r}} \frac{c}{1 + 2iQ_i \frac{\delta f}{f_r}}.$$

where 1/Q₁ = 1/Qᵣ + 1/Qₑ. Here, Qᵣ is the material and temperature dependent internal quality factor of the λ/4 structure, and Qₑ the geometry dependent coupling quality factor. Qₑ is a complex-valued parameter closely related to the coupling quality factor Qₑ, whose finite phase Θ can give rise to a line shape asymmetry due to non-ideal circuit elements, e.g. a complex loading of the resonator or impedance mismatches. Its real part fulfills the condition Re[Qₑ⁻¹] = Qₑ⁻¹. Fig. 5(b) shows the normalized data of the transmission |S₂₁|² for one exemplary resonance of a pristine film at T = 100 mK. It includes a fitted curve following Eq. 2, an intrinsic quality factor of Qᵣ = 23600 and a coupling quality factor of Qₑ = 10800 are obtained. Fig. 5(c) displays an overview plot of the uncalibrated transmission of a similar device, this time after an exposure of 30 min to the hot CVD gas mixture. Still, all three resonances can be clearly recognized, with quality factors up to Qᵣ = 5000.

VII. COPLANAR RESONATOR TEMPERATURE DEPENDENCE

Fig. 6 shows the observed temperature dependence of the resonance frequency fᵣ(T) (Figs. 6(a,c)) and the internal quality factor Qᵣ(T) (Figs. 6(b,d)) for two devices. The device of Figs. 6(a,b) (device 1) has been characterized after thin film deposition and patterning, the device of Fig. 6(c,d) (device 2) has been additionally exposed to the nanotube CVD growth environment.

As can be seen in the figure, the resonance frequency fᵣ clearly decreases at high temperatures (T ≥ 0.8 K for device 1, T ≥ 0.4 K for device 2). This can be attributed to a decrease in superfluid density of the superconducting thin film. A related rise in the quasiparticle density leads to a higher damping of the resonator and therefore to a reduction of the internal quality factor. Following Mattis and Bardeen [34, 37], the temperature dependence of fᵣ can be expressed as

$$\frac{\delta f_r}{f_0} = \frac{\sigma_0}{\sigma_2} \frac{\delta \sigma_2}{\sigma_2},$$

with δfᵣ(T) = fᵣ(T) - f₀ the deviation of the resonance frequency with finite temperature and f₀ = fᵣ(T = 0) the approximated resonance frequency at T = 0. In the limit hf ≪ Δ(T = 0) and k₆T ≪ Δ(T = 0) the imaginary part σ₂ of the complex conductivity σ of the device can be approximated using the temperature-dependent BCS energy gap Δ and the normal state conductivity σᵣ as [38]

$$\frac{\sigma_2}{\sigma_0} = \frac{\pi \Delta}{h f} \left[ 1 - 2 e^{-\Delta/k_B T} e^{-h f/2k_B T} f_0 \left( \frac{h f}{2k_B T} \right) \right].$$
The quality factor is

$$\delta \left( \frac{1}{Q_i} \right) = \alpha_0 \frac{\delta \sigma_1}{\sigma_2}$$

with the real part of the complex conductivity [38]

$$\frac{\sigma_2}{\sigma_n} = 4\Delta \frac{h}{h_f} e^{-\Delta \langle k_B T \rangle} \sinh \left( \frac{h f}{2 k_B T} \right) K_0 \left( \frac{h f}{2 k_B T} \right).$$

Here $K_0(x)$ is a modified Bessel function of the second kind. Using the parameters extracted in Figs. 6(a,c) we can plot the corresponding expected temperature dependence $Q_i(T)$, again as solid red and dashed green lines; while significant deviations exist, the overall tendency agrees well with the data.

For $T \ll T_c$, the theory by Mattis and Bardeen predicts no change in the resonance frequency with temperature. However, in our devices for $T \lesssim 0.5$ K a slight decrease in $f_r$ is observed, leading to an overall nonequilibrium behaviour of $f_r(T)$, see the insets of Figs. 6(a,c). This is consistent with the influence of two-level systems (TLS) in the substrate contributing to both dissipation and dispersion and can be described by [40–42]

$$\frac{\delta f_r}{f_0} = F \frac{\delta \varepsilon}{2} \varepsilon = F \frac{\vartheta\pi \sigma_{\vartheta}}{\pi} \left[ \Re \Psi \left( \frac{1}{2} + \frac{1}{2\pi i} \frac{h f_r(T)}{k_B T} \right) \right] + \ln \left( 1 + \frac{h f_r(T)}{k_B T} \right).$$

Here, $\Psi$ is the digamma function, $0 < F < 1$ the filling factor, giving the ratio of the electric energy stored in the TLS hosting material to the total electric energy stored in the resonator, and $\vartheta$ the loss tangent of the substrate. The insets of Figs. 6(a,c) show the measured data along with a low-temperature fit, using the product $F \vartheta$ and $f_0$ as fit parameters.

In the case of device 1, Fig. 6(a), this results in $F \vartheta = 3.968 \cdot 10^{-5}$. This value of $F \vartheta$ is comparable to literature values for niobium resonators on sapphire, see e.g. [43]. The fit value for $F \vartheta$ provides an approximation for the TLS-related low-power, low-temperature internal Q-factor $Q_i,\text{TLS} = 1/(F \vartheta) \approx 25000$. As expected this slightly exceeds our measured value of $Q_i(T = 20 \text{ mK}) \approx 23500$. The fit for device 2 provides $F \vartheta = 4.087 \cdot 10^{-5}$. This results in a quality factor $Q_i,\text{TLS} \approx 24500$, comparable to the value for device 1 and apparently insensitive to the degradation of the device due to the CVD process.

Similar to $f_0$, also the low-temperature behavior of the quality factor $Q_i$ is governed by interactions with substrate TLS. This can be modeled by [41, 42]

$$\frac{1}{Q_i} = F \vartheta_{\text{eff}} \tan \left( \frac{h f_r(T)}{2 k_B T} \right) + \frac{1}{Q_{\text{other}}}.$$  

Here, $\vartheta_{\text{eff}}$ is an effective, reduced loss tangent which takes into account that at strong driving the two-level systems are partially saturated and thereby unable to absorb energy. $Q_{\text{other}}$ describes dissipative processes unrelated to TLS. The fit for device 1 again agrees well with

where $I_0(x)$ is a modified Bessel function of the first kind. The parameter $\alpha_0$ in Eq. 3 is the kinetic inductance fraction in zero-temperature limit of the coplanar waveguide.

The solid red lines in Figs. 6(a,c) are fit curves corresponding to this model, using in each case $\alpha_0$ as a free parameter. The result agrees very well with the experimental data. For device 1 (no CVD) we obtain $\alpha_0 = 0.199$, for device 2 (after 10 min CVD) $\alpha_0 = 0.249$. The value of $\alpha_0$ can be calculated from the normal-state conductance $\sigma_n$ and the critical temperature $T_c$, following [39]. Using the parameters of our devices, we obtain $\alpha_{\text{th}} = 0.243$ and $\alpha_{\text{th}} = 0.284$, in reasonable agreement with the fit results. The corresponding functional dependence of $f_r(T)$ is plotted in Figs. 6(a,c) each as a green dashed line.

In Figs. 6(b,d), clearly also the internal quality factor $Q_i$ of the devices decreases at high temperature. Following Mattis and Bardeen, the corresponding change of the
the measured data (inset of Fig. 6(b)) and results in $F_{\text{eff}} = 5.677 \times 10^{-7}$ and $Q_{\text{other}} = 23800$. For device 2 (inset of Fig. 6(d)) only few data points at low temperature are available, and a saturation of the device temperature cannot be excluded. Applying the fit, we obtain $F_{\text{eff}} = 1.634 \times 10^{-6}$ and $Q_{\text{other}} = 4370$, dominating the TLS contribution.

As can be seen in Fig. 6, the combined effect of two-level systems and a temperature dependent superfund density provide a good description of our devices. Due to the CVD process the kinetic inductance fraction increases. The low temperature fits tentatively display similar influence of two-level systems before and after CVD, however a significant non-TLS induced dissipation term results in the second, post-CVD device.

VIII. SUMMARY AND CONCLUSIONS

The x-ray photoelectron spectroscopy characterization demonstrates that our co-sputtering process from two independent targets can generate MoRe thin films of controlled alloy composition. We observe traces of molybdenum oxides at the film surface after ambient air exposure. Exposing the thin films to the carbon nanotube CVD growth environment, a significant amount of carbon is incorporated into the film. No clear indications for the formation of molybdenum carbide can be found; the surface oxide is absent after CVD, pointing towards its reduction in the 850°C H₂ atmosphere.

In electrical characterization, we observe a higher resilience to the CVD process for MoRe$_{30}$ films; the critical temperature only drops below $T_c = 8$ K for growth times $> 20$ min. For shorter growth times, data indicates that the CVD process may have effects similar to the annealing discussed in [32], i.e., enhancing the critical temperature and critical current density. We observe up to $j_c = 2.7 \times 10^5$ A/mm$^2$ and $T_c = 9.2$ K in a Hall bar geometry.

$\lambda/4$ coplanar waveguide resonators were defined by thin film deposition and subsequent reactive ion etching. After structuring, we observe internal quality factors up to $Q_i = 37000$ in pristine devices and up to $Q_i = 5000$ after undergoing the CVD process. The temperature evolution of the resonance frequency $f_i$ and the internal quality factor $Q_i$ of the devices can be understood for low temperatures $T \leq 0.4$ K via interaction with substrate two-level systems, at higher temperatures $T \geq 0.8$ K via the decreasing superfund density.

Mo$_{20}$Re$_{80}$ clearly excels in terms of stability during the CVD process, critical temperature and field as well as critical current. However, the achieved quality factors are clearly lower than those reported in literature for similar materials [34]. Our reference niobium devices characterized for comparison have reached $Q_i = 4 \cdot 10^6$, excluding the detection setup as cause. The effect of the CVD process in lowering $Q_i$ is similar to other published observations [34]. Further improvements thus should be targeted at the radiofrequency properties of the substrate as well as the pristine, as-deposited metal film and coplanar waveguide resonators defined in it.

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