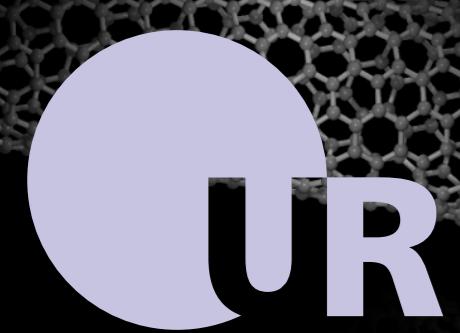


Shaping electron wave functions in a carbon nanotube with a parallel magnetic field^[1]

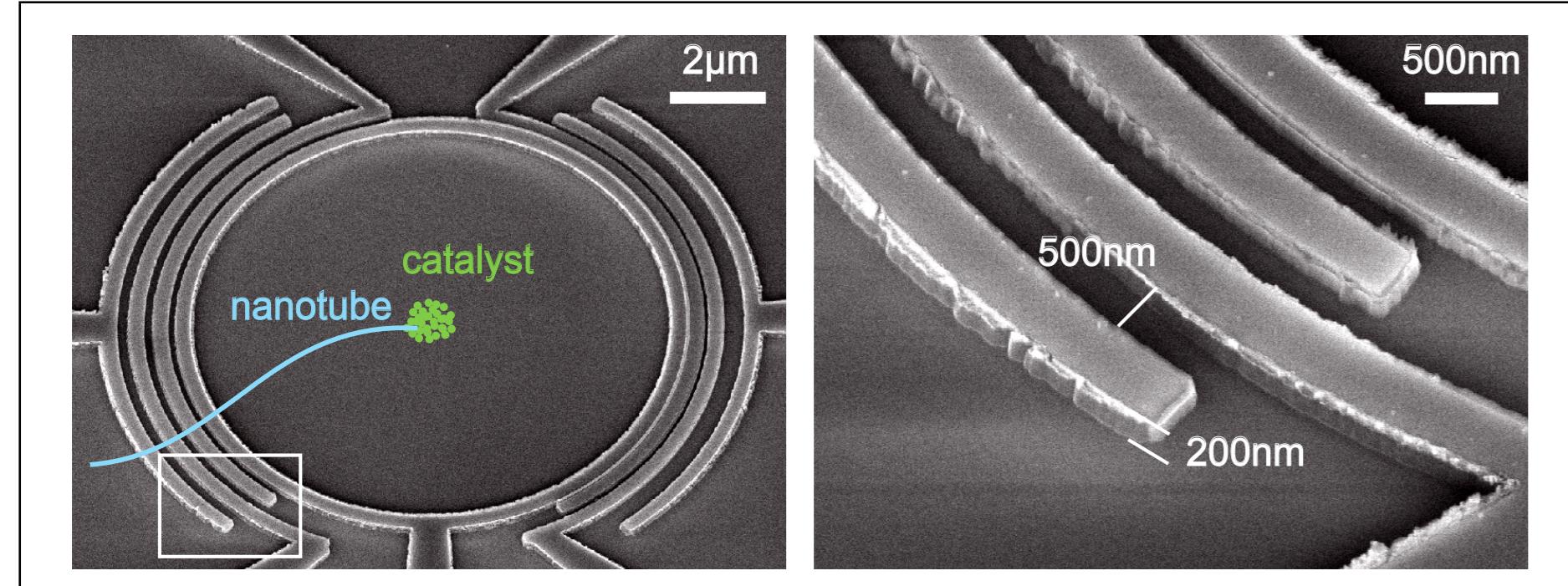


M. Margańska, D. Schmid, A. Dirnaichner, P. Stiller, Ch. Strunk, M. Grifoni, and A. K. Hüttel

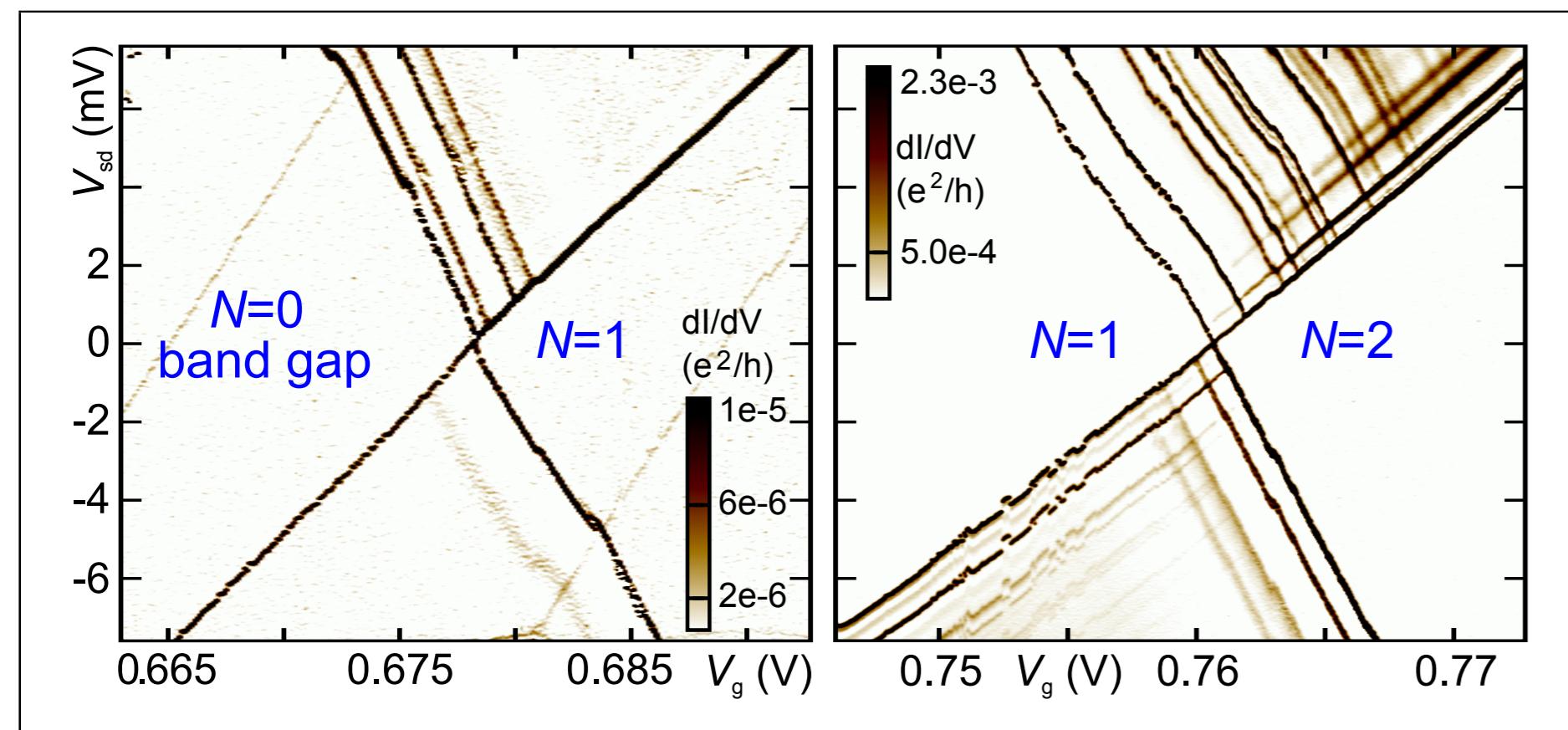
University of Regensburg, 93040 Regensburg, Germany



Ultraclean carbon nanotubes



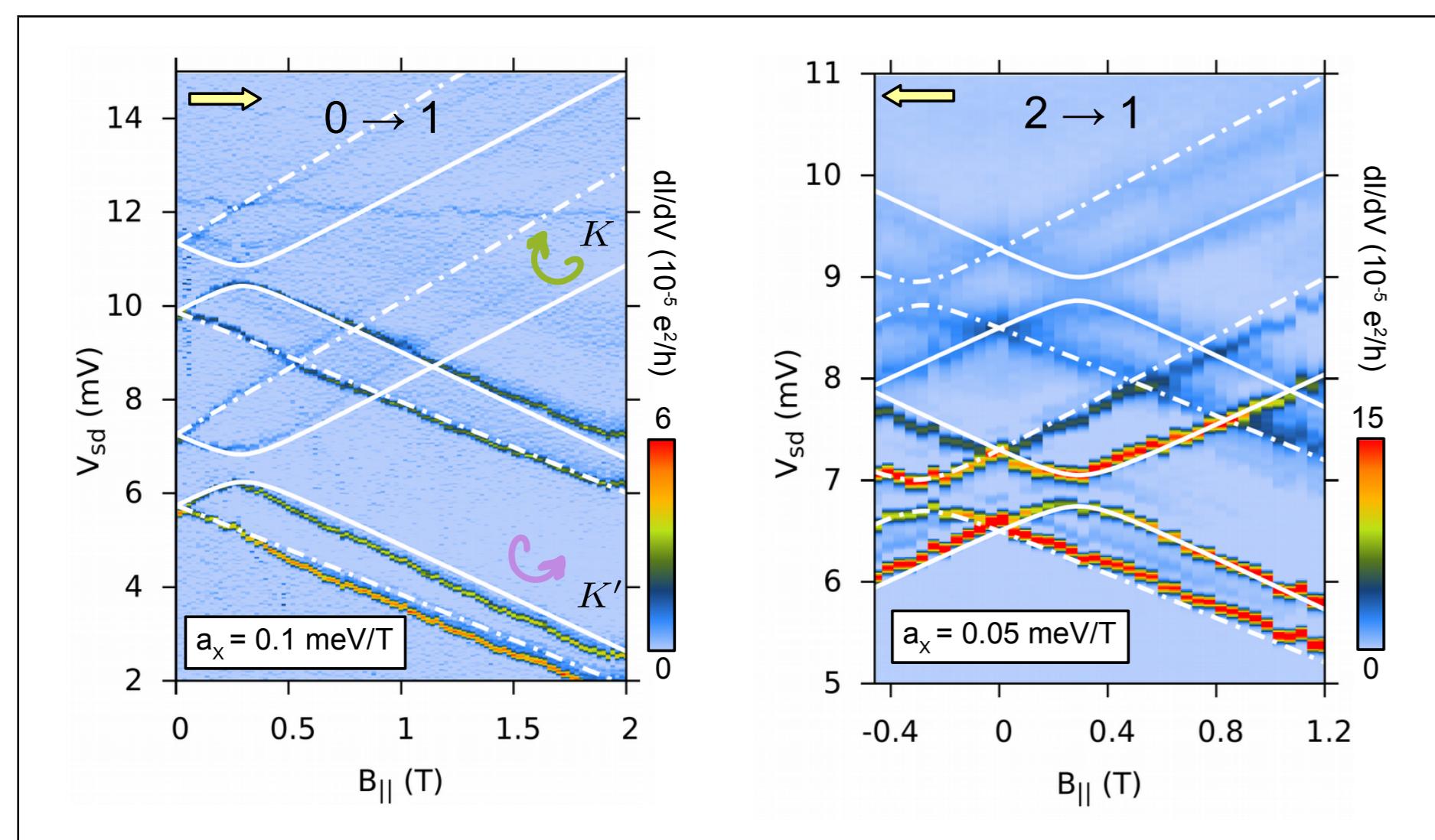
- first prepare metal structures and trenches
- then grow nanotubes across leads
- no lithography or wet chemistry afterwards
→ no chemical or mechanical damage
- no resist residues, no e-beam irradiation
- clean few-electron system [1, 2, 3, 4]
- spectroscopy of one- and two-electron states



N=1: Low field behaviour

- linearized single particle Hamiltonian

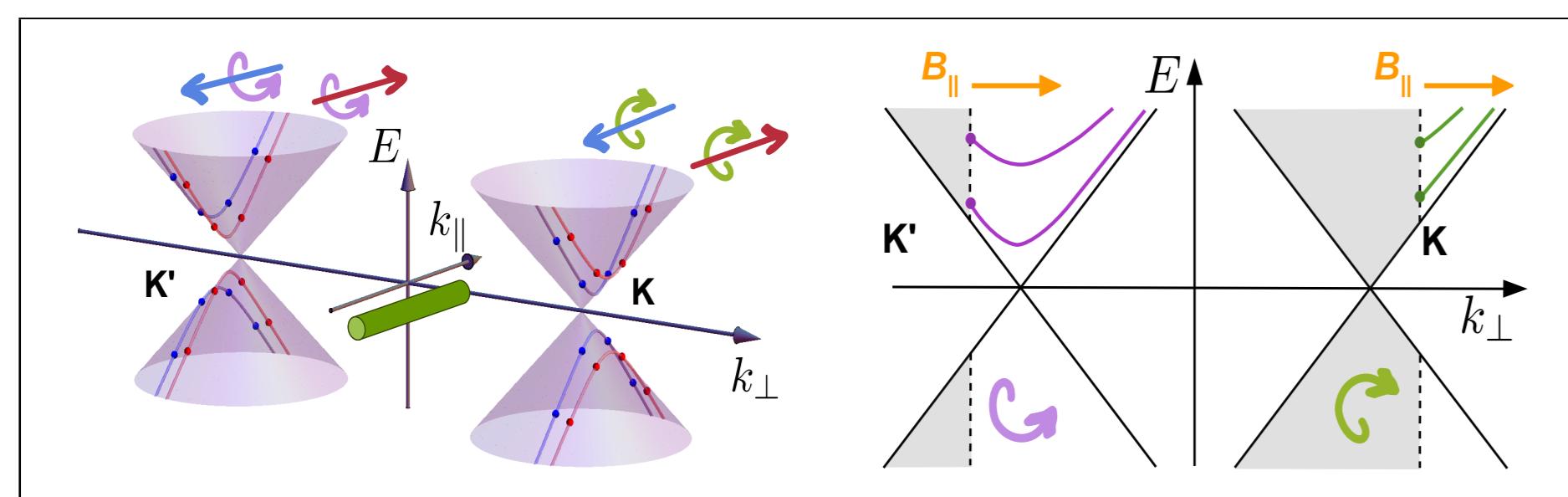
$$\hat{H}_{\text{CNT}} = \varepsilon_d \hat{I}_\sigma \otimes \hat{I}_\tau + \frac{\Delta_{KK'}}{2} \hat{I}_\sigma \otimes \hat{\tau}_z + \frac{\Delta_{SO}}{2} \hat{\sigma}_z \otimes \hat{\tau}_x + \frac{g_s \mu_B |\vec{B}|}{2} (\cos \varphi \hat{\sigma}_z + \sin \varphi \hat{\sigma}_x) \otimes \hat{I}_\tau + g_{orb} \mu_B |\vec{B}| \cos \varphi \hat{I}_\sigma \otimes \hat{\tau}_x$$



- extract high-level parameters Δ_{SO} , $\Delta_{KK'}$, $\Delta_{||}$, ...
- fit requires additional corrections

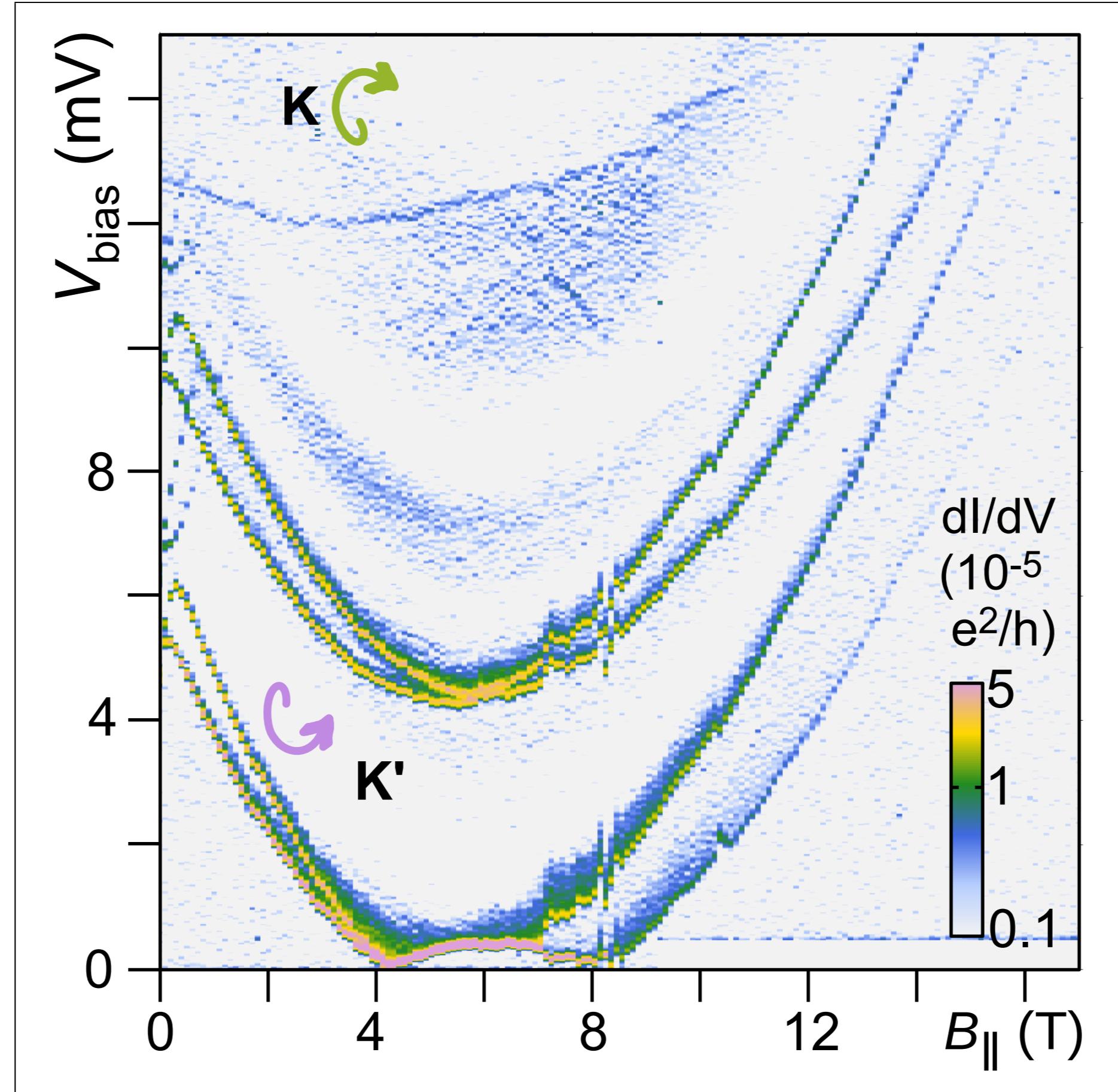
Full dispersion relation [5]

- $B_{||}$ "scans" Dirac cone



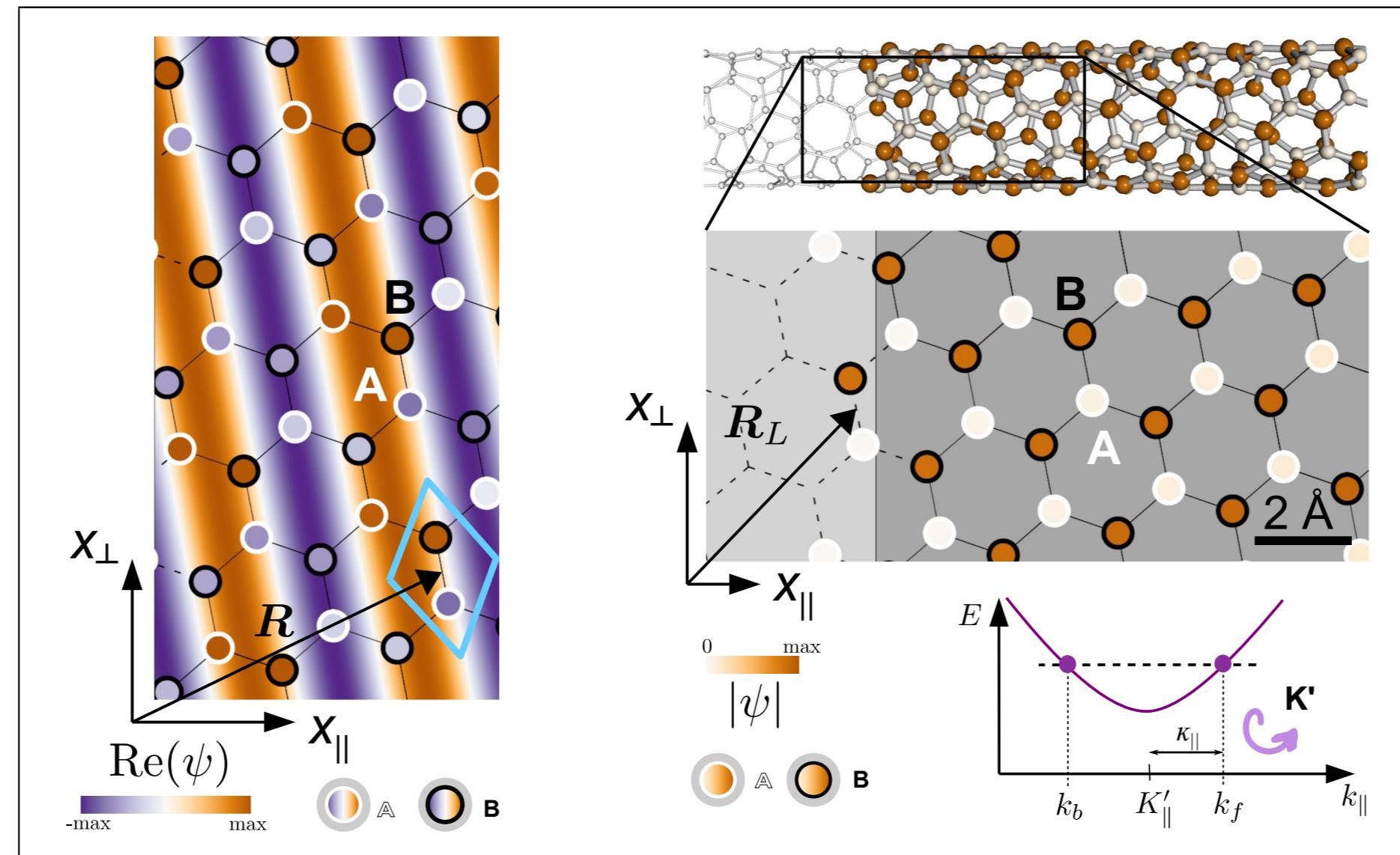
$$E_\pm(\tau, \sigma, \kappa_{||}, B_{||}) = \tau \sigma \varepsilon_{SO} + \sigma \mu_B B_{||} + \hbar v_F \sqrt{\left(\tau \Delta k_{\perp} + \sigma \Delta k_{SO} + \frac{\pi R}{\phi_0} B_{||} \right)^2 + \kappa_{||}^2}$$

N=1: High field behaviour



- K conductance peaks fade out very quickly
- K' conductance increases, then decreases

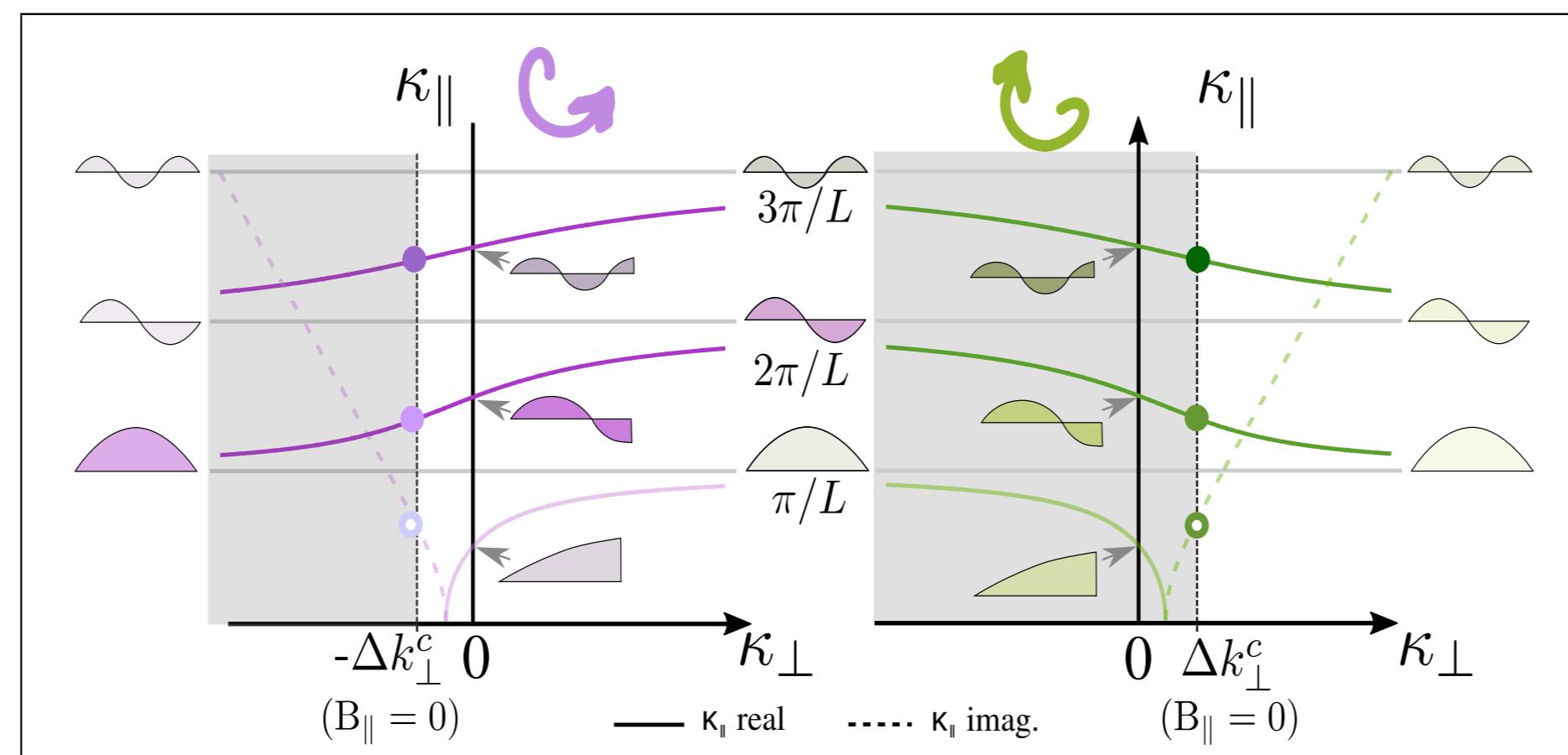
Boundary conditions



- bipartite lattice → different phase on the two sublattices (advanced / retarded)
- construct standing waves (forward + backward)
- no superposition makes wavefunction vanish on both sublattices at both ends
- minimize dangling bonds → one sublattice has majority at one end, other at the other end
- not a quantum "box", more a " $\lambda/4$ resonator"; "cross-quantization" of κ_{\perp} and $\kappa_{||}$

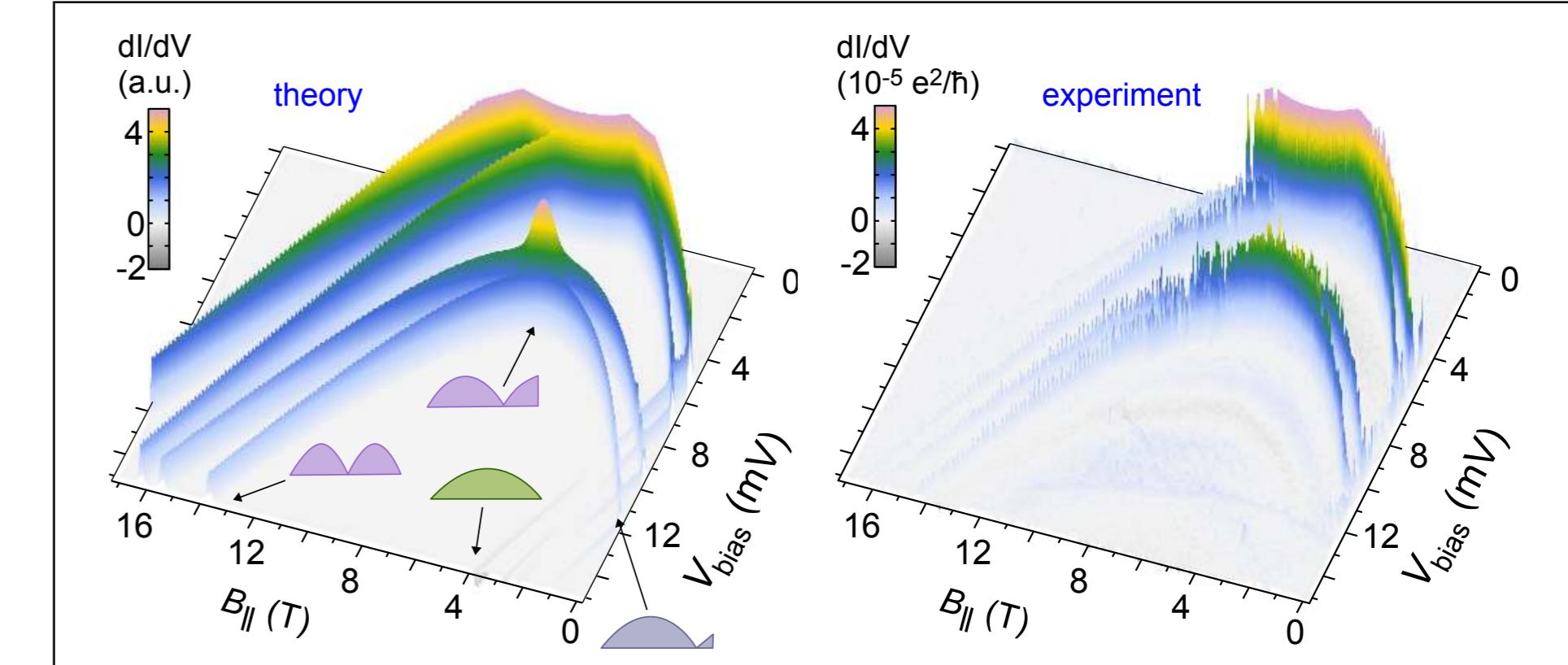
$$e^{2i\kappa_L} = \frac{\tau \kappa_{\perp} + i \kappa_{||}}{\tau \kappa_{\perp} - i \kappa_{||}}$$

Wave function shapes in $B_{||}$



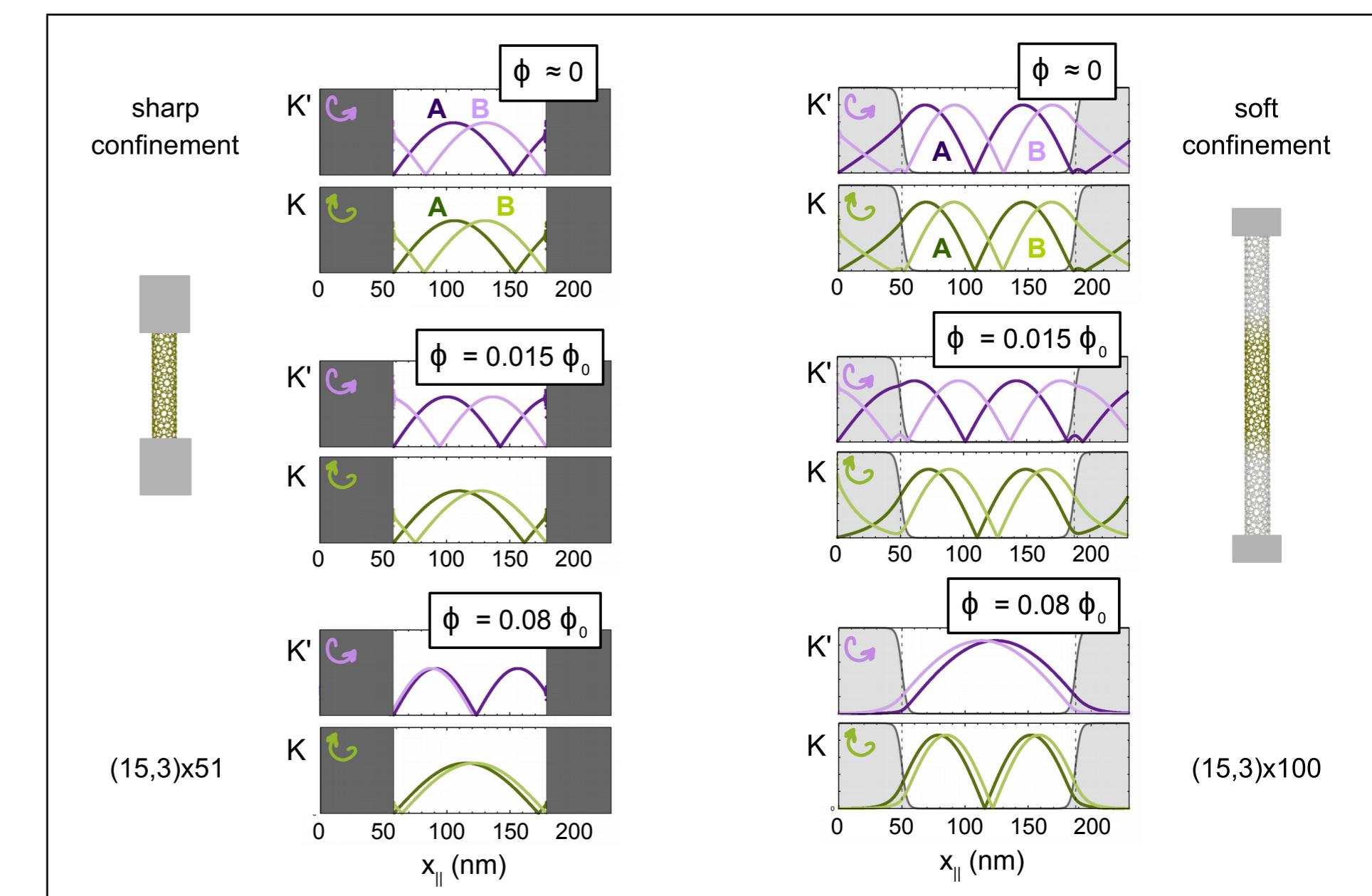
- magnetic field $B_{||}$ selects κ_{\perp} (vertical dotted line)
- field shifts nodes of wave function to the end; " $\lambda/4$ to $\lambda/2$ " tuning

Impact on conductance



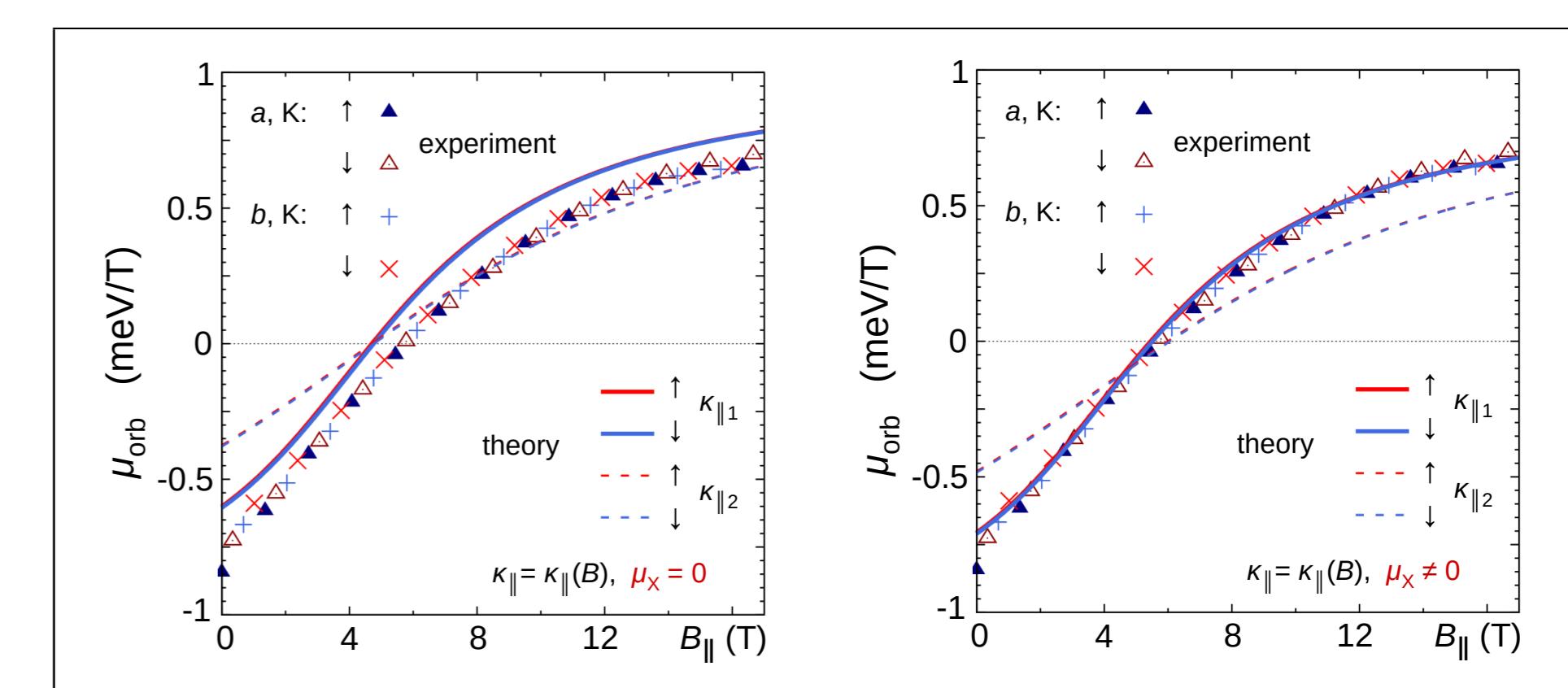
- tunnel coupling ↔ wave function amplitude at the nanotube ends
- introduce into model → very good agreement
- (initially) downsloping lines, ↗: from $\sim \lambda/2$ shape to $\lambda/4$, then to $\lambda/2$; dI/dV increases, then decreases again
- upsloping lines, ↘: dI/dV decreases fast, wavefunction approaches $\lambda/2$ shape

Soft confinement potential

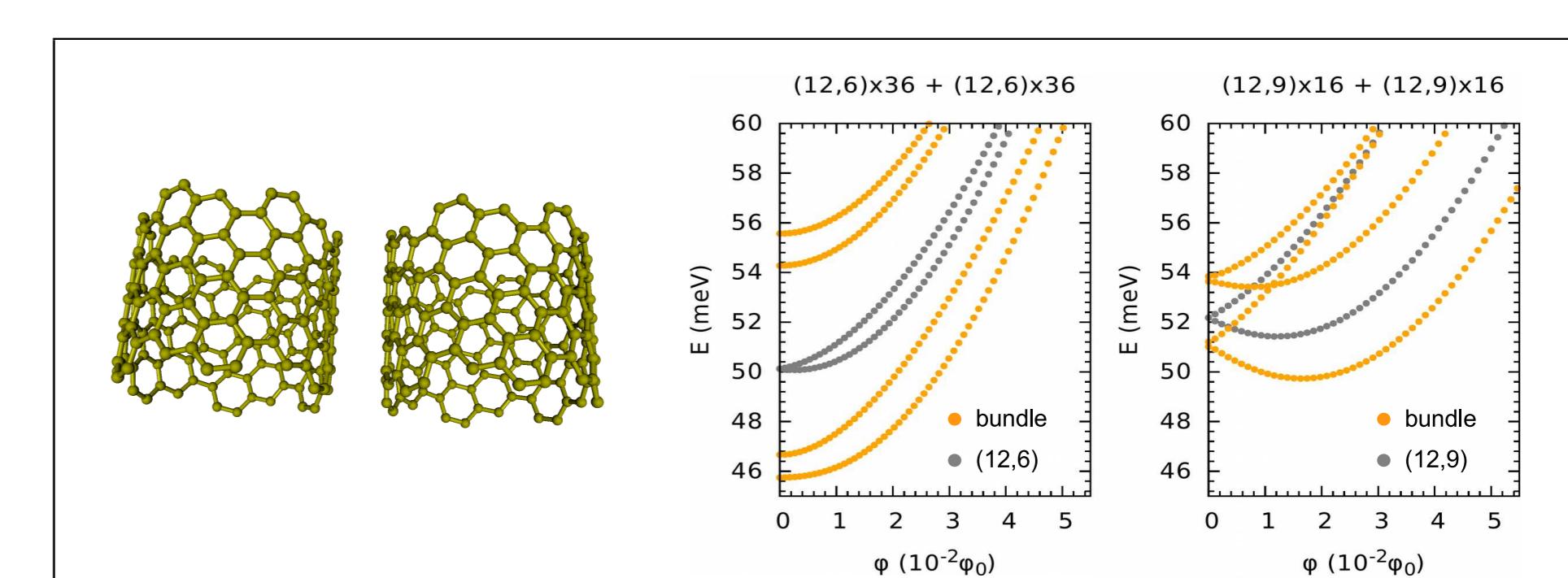


- qualitatively same result for wave functions

Origin of the second shell?



- anomalous magnetic moment, "offset"?
- low-lying (1.6 meV) second shell: same $\kappa_{||}$ (!)
- additional degeneracy; bundle of two CNTs?



References

- [1] M. Margańska *et al.*, arXiv:1712.08545 (2018).
- [2] D. R. Schmid *et al.*, PRB **91**, 155435 (2015).
- [3] A. Dirnaichner *et al.*, PRL **117**, 166804 (2016).
- [4] K. J. G. Götz *et al.*, PRL **120**, 246802 (2018).
- [5] W. Izumida *et al.*, JPSJ **78**, 074707 (2009).
- [6] S. Reinhardt *et al.*, arXiv:1804.03321 (2018).